

Crystallographic, magnetic and calorimetric studies of $\text{Ho}_5\text{Si}_2\text{Ge}_2$

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Abstract

Following the discovery of a giant magnetocaloric effect in $\text{Gd}_5(\text{Si,Ge})_4$, attention has been extended to $\text{R}_5(\text{Si,Ge})_4$ with R being other rare-earth elements. In this work, X-ray structural analyses, low- and high-field magnetization measurements and zero-field calorimetric measurements were carried out on $\text{Ho}_5\text{Si}_2\text{Ge}_2$. Specific heat data were also obtained for nonmagnetic $\text{Lu}_5\text{Si}_2\text{Ge}_2$ as a reference material. In contrast to the general trend of having a ferromagnetic order in the $\text{R}_5(\text{Si,Ge})_4$ series, $\text{Ho}_5\text{Si}_2\text{Ge}_2$ actually becomes antiferromagnetically ordered with a Néel temperature T_N near 25 K. Moreover, an anomalous behavior below T_N also prevails in the temperature dependence of both magnetization and specific heat, suggesting further transitions from the antiferromagnetic to other complex magnetic structures.

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1. Introduction

Since the discovery of a giant magnetocaloric effect (GMCE) in $\text{Gd}_5\text{Si}_2\text{Ge}_2$ [1,2], this compound has taken the central stage of research activities on magnetic refrigeration materials. A preliminary study on the $\text{R}_5\text{Si}_2\text{Ge}_2$ series with other rare earths

(R) completely replacing Gd revealed quite interesting magnetic ordering phenomena. It has been shown that, similar to $\text{Gd}_5(\text{Si,Ge})_4$, Tb-based compounds also undergo ferromagnetic transitions with T_c near room temperature, along with a GMCE in its vicinity [3,4]. Meanwhile, other R-based compounds order at lower temperatures and show complex magnetic structures [3]. As an extension of this line of research, we report here on crystallographic, magnetic and calorimetric studies of $\text{Ho}_5\text{Si}_2\text{Ge}_2$. To delineate the thermal

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property into lattice and magnetic contributions, calorimetric measurements were also made on $\text{Lu}_5\text{Si}_2\text{Ge}_2$ as a nonmagnetic reference.

2. Experimental

The $\text{Ho}_5\text{Si}_2\text{Ge}_2$ and $\text{Lu}_5\text{Si}_2\text{Ge}_2$ samples were prepared by arc-melting a stoichiometric mixture of Ho/Lu (3 N), Si (5 N) and Ge (4 N) in a pure Ar atmosphere. The resulting ingots were turned over and remelted several times to ensure sample homogeneity. The overall weight loss was less than 0.6%. Whereas the as-melted $\text{Lu}_5\text{Si}_2\text{Ge}_2$ ingot was directly used for measurements, the $\text{Ho}_5\text{Si}_2\text{Ge}_2$ ingot was further sealed in a quartz ampoule with pure Ar atmosphere and annealed for 7 days at 1000°C , followed by water quenching. The crystallographic structure and the sample quality were studied using X-ray diffraction and electron probe microanalysis (EPMA). The magnetic properties of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ between 2 and 300 K were determined with a 5 T SQUID magnetometer. Further magnetization measurements in high fields up to 40 T were carried out at the Amsterdam High Field Installation. The specific heat of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ and $\text{Lu}_5\text{Si}_2\text{Ge}_2$ was measured in the temperature range of 0.6–40 K by using a microcalorimeter as described elsewhere [5].

3. Results and discussion

The room-temperature XRD-patterns of the $\text{Ho}_5\text{Si}_2\text{Ge}_2$ and $\text{Lu}_5\text{Si}_2\text{Ge}_2$ samples are presented in Fig. 1. For $\text{Ho}_5\text{Si}_2\text{Ge}_2$, the analysis indicates a single phase of the orthorhombic structure, space group Pnma , with unit cell parameters of $a = 7.501 \text{ \AA}$, $b = 14.510 \text{ \AA}$, and $c = 7.599 \text{ \AA}$. In reconfirming this by the EPMA analysis, the matrix of the EPMA image in the inset of Fig. 1 corresponds to the main phase having an exact composition of $\text{Ho}_5\text{Si}_{1.6}\text{Ge}_{2.1}$. The dark phase in the image represents an “impurity” phase of the composition $\text{Ho}_5\text{Si}_{2.8}\text{Ge}_{1.6}$, which amounts to only a few percent in volume. For $\text{Lu}_5\text{Si}_2\text{Ge}_2$, an orthorhombic structure was determined with

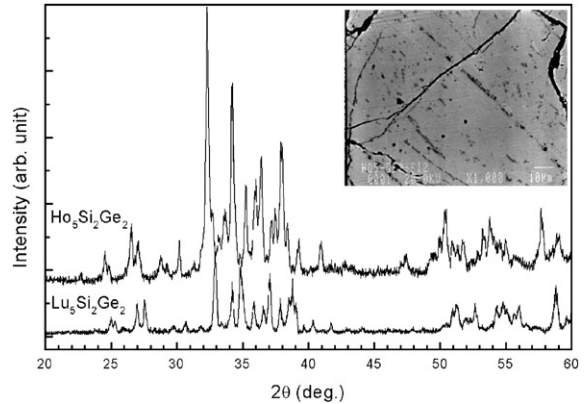


Fig. 1. Room-temperature XRD-patterns of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ and $\text{Lu}_5\text{Si}_2\text{Ge}_2$ powders. The inset shows an EPMA image of the $\text{Ho}_5\text{Si}_2\text{Ge}_2$ sample.

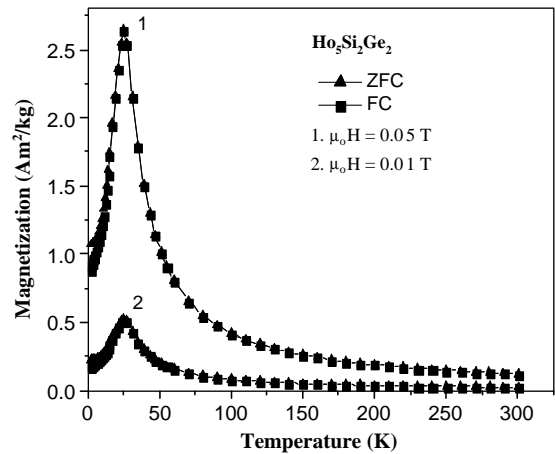


Fig. 2. Temperature dependence of FC- and ZFC-magnetization of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ in a field of 0.01 and 0.05 T.

lattice constants $a = 7.386 \text{ \AA}$, $b = 14.262 \text{ \AA}$ and $c = 7.456 \text{ \AA}$.

The temperature dependence of the magnetization of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ is given in Fig. 2. A magnetic order sets in at a much lower temperature than that in $\text{Gd}_5\text{Si}_2\text{Ge}_2$ [1]. Judging from the occurrence of a peak in the figure, the transition is of an antiferromagnetic type. The peak corresponding to the Néel temperature, $T_N \approx 25 \text{ K}$, is nearly unchanged with increasing applied field. The susceptibility shown in Fig. 3 follows a Curie–Weiss relation $\chi = C/(T - \theta)$ with a paramagnetic Curie temperature $\theta = 17 \text{ K}$. The effective Ho-

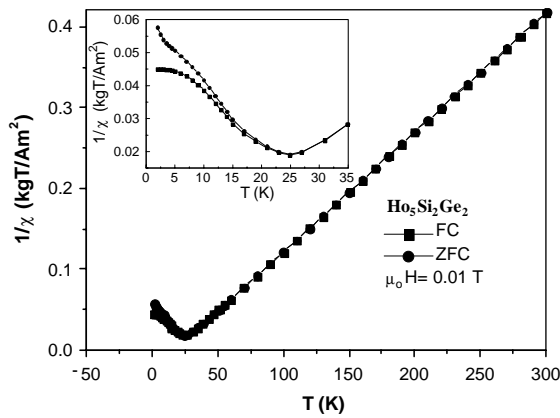


Fig. 3. Temperature dependence of the inverse susceptibility, as calculated from the magnetization data in Fig. 2, of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ in a field of 0.01 T. The inset shows a split between FC- and ZFC-data at the low-temperature region.

moment derived from these data is $9.7\mu_{\text{B}}$, compared with $g_J[J(J+1)]^{1/2}\mu_{\text{B}} = 10.60\mu_{\text{B}}$ for free Ho^{3+} . Also revealed in the inset of Fig. 3 is a split of the temperature dependence of the inverse susceptibility between the zero-field-cooled (ZFC) and field-cooled (FC) curves at $T = 15$ K, as well as an anomaly near 2 K.

Below the Néel temperature, magnetization data at different temperatures in low fields up to 5 T in Fig. 4 exhibit a metamagnetic transition. As can be seen in the inset showing the field dependence of the susceptibility, the critical field decreases with increasing temperature. The metamagnetic transition might be related to a field-induced transformation from the antiferromagnetic to some other magnetic configurations yet to be identified. This appears to be consistent with the magnetization data at 4.2 K in high fields up to 38 T in Fig. 5. They approach saturation very slowly, suggesting a quite large magneto-crystalline anisotropy and/or a complex ordered spin structure at low temperatures. By extrapolating the magnetization data to an infinite field based on $M = M_s + aH^{-1}$, the saturation magnetization moment at 4.2 K is estimated to be $10.18\mu_{\text{B}}$ per Ho^{3+} ion. This is in good agreement with the theoretical value of $g_JJ = 10\mu_{\text{B}}$. Data not shown here from magnetic hysteresis loop measurements at different temperatures below T_{N} gave negligible coercivity values.

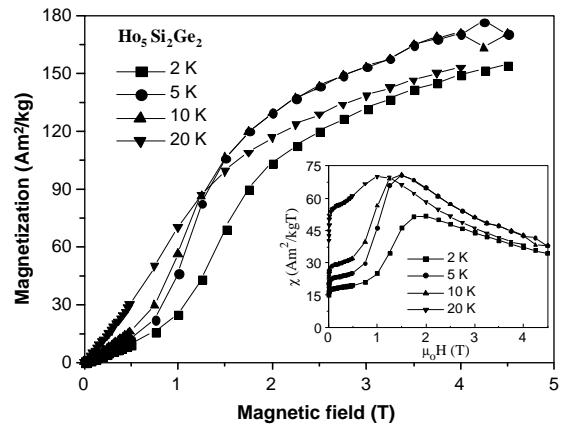


Fig. 4. Low-field magnetization of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ as a function of field at different temperatures below T_{N} . The inset shows the field dependence of the susceptibility as extracted from the corresponding magnetization curves.

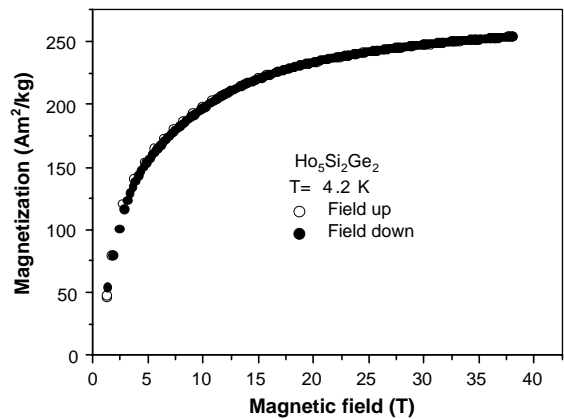


Fig. 5. High-field magnetization of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ as a function of field at 4.2 K.

Figs. 6 and 7 present specific heat data in terms of the temperature dependence of C and C/T for $\text{Ho}_5\text{Si}_2\text{Ge}_2$ and $\text{Lu}_5\text{Si}_2\text{Ge}_2$, respectively. The smooth and monotonically increasing values for nonmagnetic $\text{Lu}_5\text{Si}_2\text{Ge}_2$ in Fig. 7 are fitted to a polynomial function, which is then assumed to represent the lattice component to the measured specific heat of $\text{Ho}_5\text{Si}_2\text{Ge}_2$. The difference of these plots represents the expected magnetic contributions, at least for temperatures above 5 K, below which C/T rises steeply as part of a nuclear Schottky term as observed in holmium [6] and Ho-based compounds [7].

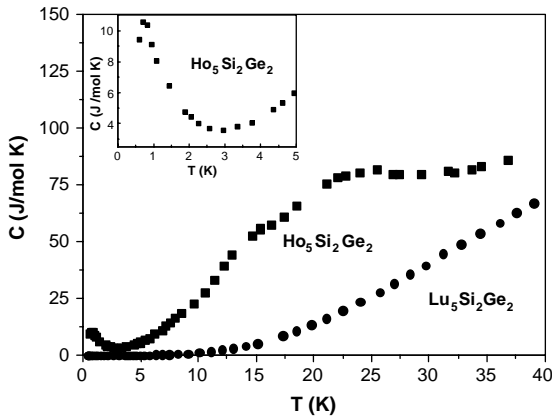


Fig. 6. Temperature dependence of the specific heat of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ and $\text{Lu}_5\text{Si}_2\text{Ge}_2$. The inset shows the lowest temperature part of the $C(T)$ curve of $\text{Ho}_5\text{Si}_2\text{Ge}_2$.

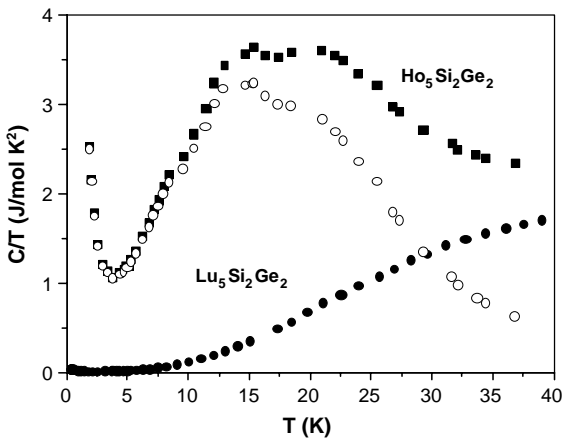


Fig. 7. Temperature dependence of C/T of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ and $\text{Lu}_5\text{Si}_2\text{Ge}_2$. The difference plot (open circles) represents mainly the magnetic contributions above 5 K and a nuclear Schottky term at the lower temperatures for $\text{Ho}_5\text{Si}_2\text{Ge}_2$.

There appear to be two identifiable magnetic anomalies. In conjunction with the magnetic data in Fig. 2, the broad C/T -peak near 20 K is obviously caused by the antiferromagnetic ordering. For the second peak around 15 K, a clue can be obtained from Fig. 8, which plots the magnetic entropy of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ above 5 K as derived by area integration of the calorimetric data in Fig. 7, $S_m(T) - S_m(5\text{K}) = \int (C_m/T) dT$. By neglecting the expectedly small value of

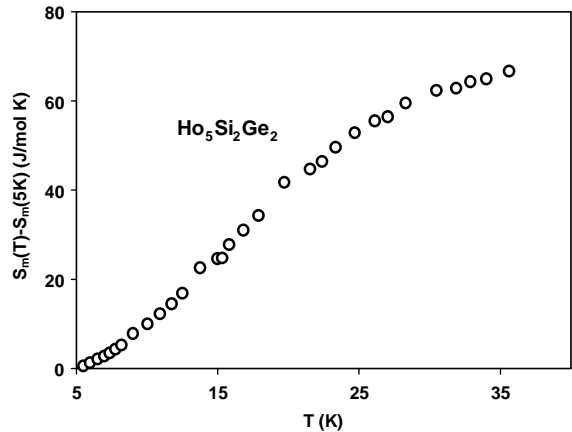


Fig. 8. Temperature dependence of the magnetic entropy of $\text{Ho}_5\text{Si}_2\text{Ge}_2$ above 5 K.

$S_m(5\text{K})$, the magnetic entropy for $\text{Ho}_5\text{Si}_2\text{Ge}_2$ reaches over 60 J/mol K above T_N , where the magnetic ordering could have vanished completely. In comparison, a magnetic entropy of only $5R \ln 3 = 46\text{ J/mol K}$ is expected for a simple order–disorder process, assuming a ground state triplet of Ho^{3+} ions. Under this consideration, the nature of the second anomaly at 15 K is difficult to establish. The entropy change points in a direction of more states involved than only the triplet ground state, complicating the magnetic phase diagram.

In conclusion, both magnetic and calorimetric measurements on $\text{Ho}_5\text{Si}_2\text{Ge}_2$ reveal an antiferromagnetic ordering at 25 K, followed by a second phase transition at 15 K of so far unknown nature. Further studies such as neutron diffraction and specific heat measurements in magnetic field are necessary to clarify the magnetic structure of both phases below 25 K.

Acknowledgements

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